Please find below our response to the comments. We thank the reviewer for careful reading of the manuscript. We address the points raised by the referee below.

The work by Rodin and Abergel revisits some of the early developments of spin relaxation theory in NMR and discusses their relation to the Lindblad equation, which is commonly employed in quantum optics and the mathematical physics literature. A detailed discussion on classical and quantum mechanical correlation functions/spectral densities with respect to detailed balance is also included.

1) The paper seems to be in response to a recent publication "A master equation for spin systems far from equilibrium" (doi.org/10.1016/j.jmr.2019.106645) authored by Malcolm Levitt and myself. However, several references and conclusions regarding our work are incorrect or misstated. In their abstract Rodin and Abergel claim that our work has "recently questioned" the "underlying theory" worked out by the founding fathers of spin relaxation theory. This is incorrect, we clearly state at the beginning of the article that our work explores the disadvantage and advantages of several thermalization techniques (IME, Jeener, Levitt-diBari) in the context of highly polarised spin systems (see section 4). In particular we show that the inhomogeneous master equation (IME) may lead to non-physical results for such systems. At no point do we cast doubt upon the underlying theory developed by Bloch, Wangness, Redfield, etc. On the contrary, we give full credit to the "founding fathers" already being "fully aware that semiclassical relaxation theory [...] is of limited validity" for highly polarised systems. If anything we cast doubt upon the current use of the IME for highly polarised spin systems within the NMR community.

We thank the reviewer for these comments, and as we alredy stated in our reply to Malcolm Levitt's « community comment » on our manuscript, we realize that the wording used in our manuscript may have been unfortunate, when it comes to allude to your paper. We would therefore want to make it clear that your insightful « Bengs & Levitt » paper points to the exact limitations of the semiclassical theory of relaxation (« high spin order », to be short). Although the assumptions that lead from the full quantum theory to this semiclassical formulation are applicable in many practical situations, the breakdown of this approximate theory and the way to overcome such limitations by using the Lindblad approach was clearly illustrated by several examples in your paper.
In our manuscript, we mention that your paper shows that the « formulation currently used by NMR spectroscopists », leads to erroneous predictions. This, of course, implicitly refers to the semiclassical theory. I hope this will dissipate any misunderstanding.
We thank the reviewer for a careful summary of his own paper. We have now modified the manuscript so as to make this point clearer.

2) Within this context section 2.2 by Rodin and Abergel seems to stipulate that we imply that it is impossible to derive the Lindblad equation from the quantized theory of Bloch and Redfield. However, we simply presented the Lindblad method as a convenient approach to extend the semiclassical relaxation theory to systems outside the weak-order limit. We rather hoped that it would be clear that the Lindblad equation follows from the quantized BWR theory given that the starting point of our derivation in section 5 coincides with the starting point of BWR. However, it is true that we could have stated this fact in more detail as has recently been done by Tom Barbara (https://doi.org/10.5194/mr-2-689-2021).

Reviewer #1 seems to overinterpret, or misinterpret our intentions, which we obviously need to clarify. In our manuscript, we wanted to show how the properties of the correlation functions emerge from the derivation of the master equation and how they relate to a semi-classical relaxation theory. Such properties are the consequences of both the Boltzmann equilibrium assumption for the lattice and of the non-commutation of the bath operators. Indeed, the Lindblad relaxation operator appears naturally, together with the associated spectral density functions that bear the important

asymmetry properties, from the perturbative solution of the evolution equation. These spectral properties and the associated Kubo relationships are direct consequences of the assumptions made on the lattice and lattice operators. Note in passing that the Kubo relations were originally defined for symmetrized correlation functions [J. Phys. Soc. Jpn (1957) 12, 570], whereas from the derivation of the master equation there emerge unsymmetrized correlation functions and spectral density functions (termed « left » and « right » in our manuscript).

The Lindblad equation makes no such assumption, and only assumes that spectral density functions reflect a Markovian process. It therefore requires additional adjustments to directly lead to a semiclassical theory. Indeed, because the bath operators B^q do not commute, the correlation functions do not obey the general symmetry rules of classical correlation functions (see our Eq. 50). However, symmetrized correlation functions do, and so are introduced (Eq. 52). It is therefore the quantum mechanical master equation with these symmetrized (quantum-mechanical) correlation functions that should be used for a semi-classical theory (which we call « pseudo-classical » to distinguish it from the theory where the effect of the bath is taken into account only through random functions). We here recovered a result obtained by Hubbard.

Finally, we may have indeed misinterpreted misleading section titles in your paper, where paragraphs 4 and 5 are explicitly devoted to 'thermalization' issues, which apparently wrongly led us to believe that this was the focus. Incidently, we do not believe that « Lindblad thermalization » is adequate, as the purpose of the Lindblad approach is to provide a general form of Markovian quantum dissipative systems. There is no assumption as to the equilibrium state, the latter being assigned a particular form based on the additional, thermodynamical, hypothesis of the bath (lattice) being in a Boltzmann equilibrium. In this sense, there are more assumptions in the NMR relaxation theory than in the Lindblad approach, and the Bloch theory yields both the Kubo relations and the Lindblad equation in the special case of spin relaxation.

We have modified the manuscript so as to make this point clearer.

3) The authors conclude section 2.2 with "the Lindblad equation [...] is thus derived from the usual quantized theory of relaxation." I am not sure if the authors are trying to claim novelty for this result (it also appears to be core statement of their abstract), if so they should consult the books by Breuer (DOI:10.1093/acprof:oso/9780199213900.001.0001), Schaller (10.1007/978-3-319-03877-3), and the recent article by Tom Barbara.

We have assumed that this remark from reviewer #1 was a purely rethorical one, as our manuscript (including the title) insistently tends to demonstrate that, as Tom Barbara put it in a comment, « Bloch got it right from the beginning ».

4) Unfortunately some mischaracterisation of our analysis of singlet-triplet relaxation in highly polarised spin systems may be found within the conclusion. The authors state: "The author's interpretation was that "thermalization" was not properly taken into account by the inhomogeneous semiclassical version of the Redfield equation, which was ascribed to the fact that the classical approach does not comply with the detailed balance condition." No such statements are made. The deviation between these results was not ascribed to detailed balance, but to the fact that the IME generates the unphysical situation in which singlet order may directly flow into Zeeman magnetisation (clearly illustrated in figure 1).

We agree that this may have been an inaccurate interpretation of your paper. However, this is an interesting point that requires clarification, as the reviewer's response indicates. First, the term « classical » was meant to refer to the « conventional », semi-classical, approach, and indeed appears to be a bad choice.

The question is why the IME provides erroneous predictions in some cases. From the quantummechanical approach, it is obtained in the « low order » case, when the thermal corrections to the double commutator part of the relaxation operator are retained to first order in the largest eigenvalue $\mathfrak{A}_k/kT$, and the relaxation operator reduces to a double commutator. However, the equilibrium density operator is not a stationary solution in this case and therefore a correction term is added to the master equation and gives the same result as the usual semiclassical master equation. And so, when the low order assumption is not verified, as in the case of a spin system prepared in the singlet state, this description becomes inconsistent.

The non commutation of the bath operators have critical consequences, leading to the lattice temperature dependent terms in the master equation. Only when the bath operators $B^{-q}(t)$ and $B^q(0)$ commute, one recovers the double commutator expression, with the additional property $J_L(\mathfrak{A})=J_R(\mathfrak{A})$, so that the Kubo relation imposes an infinite lattice temperature. This illustrates how the finite temperature of the lattice is conveyed to the spins through non commutation of the bath operators of the coupling hamiltonian.

In the conventional semiclassical approach, the spin-bath interactions are represented by spin hamiltonians modulated by random functions of the lattice. Therefore, it has two simultaneous consequences : the structure of the relaxation operator is affected in such a way that the master equation takes the form of a double commutator ; and since $J_L(\mathfrak{A})=J_R(\mathfrak{A})$, the system cannot evolve to a thermodynamic equilibrium associated with a finite temperature. In this case, detailed balance property is conserved but only in the special case of infinite lattice temperature.

Thus, the ad hoc equilibrium density operator term used to account for a finite lattice temperature in the semi-classical theory, which makes the master equation inhomogeneous, can reflect the dynamics of the system only in situations where the non commuting terms can be neglected, which correspond, as stated by Redfield, to a situation where the spins are prepared in « an usual way ».

The manuscript has been modified to clarify this point.

5) The provided discussion by the authors on this subject is also somewhat confusing. If I understand correctly the authors strongly emphasise the validity of detailed balance for the bath (in example the equilibrium distribution of the liquid follows Boltzmann). But this does not matter very much if the spin transitions do not obey detailed balance, which is clearly what our work refers to when talking about detailed balance. Consider for example equation (77) with a spectral density $J_\theta(\omega)=\theta(\omega)J(\omega)$, where $\theta(\omega)$ enforces detailed balance of the spectral densities. Consider now an ensemble of spin-1/2 particles experiencing random field relaxation in solution. According to equation (77) the relaxation superoperator is given by

$$\hat{\Gamma} = -\frac{1}{2}\xi_{\mathrm{erf}}\sum_{m=-1}^{+1}J(|m|\omega)\theta(m\omega)\hat{T}_{1m}^\dagger\hat{T}_{1m}$$

It is straightforward to verify the following matrix elements

$$W_{\alpha\leftarrow\beta} = (P_\alpha|\hat{\Gamma}|P_\beta) = \frac{\xi_{\mathrm{erf}}}{4}J(\omega)(\theta(+\omega)+\theta(-\omega))$$

$$W_{\beta\leftarrow\alpha} = (P_\beta|\hat{\Gamma}|P_\alpha) = \frac{\xi_{\mathrm{erf}}}{4}J(\omega)(\theta(+\omega)+\theta(-\omega)),$$

where P_α and P_β are the α and β population operators, respectively. The two transition rates per unit time $W_{\alpha\leftarrow\beta}$ and $W_{\beta\leftarrow\alpha}$ are equal. So even if the detailed balance condition holds for the bath it does not hold for the spin-state transitions when equation (77) is used. In other words it is correct to say that the equilibrium density operator for the spin system is not in the nullspace of Γ because it does not obey detailed balance for spin transitions. This has been discussed extensively in appendix A.2 of our work.

This argument does not seem totally clear to us. You seem to be willing to demonstrate that the semiclassical theory of relaxation does not lead to detailed balance for the spins, and that there is no

connection between spin and bath detailed balance. My understanding is that the « eq. 77 » invoked here is the semi-classical relaxation operator. Quite obviously, one should not expect any connection between spin and bath thermal properties, as they are not taken into account. In addition, the non commutation of bath operators of the spin-bath coupling, which is central in the quantummechanical theory, as discussed above, is irrelevant in this case.

In contrast, we should point out that, in the full quantum theory, detailed balance for the spins is indeed conveyed from the detailed balance property of the bath (under the hypothesis of a bath in Boltzmann equilibrium).

6) *Besides the very unfortunate misinterpretation of our work, the provided discussion regarding classical and quantum mechanical correlation functions is interesting and is used to introduce a "thermalized double-commutator" in section 4. The thermalized double commutator renders the relaxation superoperator Redfield like and it is shown that this approach is fully equivalent to the Lindblad method. However advantages and disadvantages, if there are any, are not discussed. For example, the thermalized double-commutator does not represent a double-commutator in a strict mathematical sense and might be awkward to incorporate into existing simulation software.*

We would agree that the term « thermalized double commutator » may lack elegance, but it reflects the underlying considerations developed in our manuscript and in this letter. The spectral density functions involved obey general symmetry properties of classical correlation functions , which can therefore be used in a straigthforward manner to generate semi-classical approximations of the master equation.

7) *Additionally, representing the relaxation superoperator in Lindblad form does not only provide a simple check on positivity and Markov compatibility, but is also in agreement with the majority of the literature outside NMR.*

The formulation used in our manuscript is totally in line with Lindblad. This has been explicitly indicated in several instances (see paragraph 3 of our manuscript, in particular.). The point that seemed important to make regarding the Lindblad formulation is that, since it says nothing about the thermodynamic equilibrium of the bath, such assumptions must be introduced explicitly in a way that is not directly connected to the physical problem at hand, in contrast to the Bloch-Redfield approach, from which critical properties of the spin correlation functions can be deduced.

8) *Apart from that there also some technical issues.*

• *The trace operator Tr should not be italicized.*
This has been corrected
• *The interaction frame transformation above equation (4) should read as* $\exp\left(-L\square_0 t\right)$
This has been corrected
• *The authors switch to a time dierence variable* $\tau = t - t'$ *in equation (6), but this is not consistent with equations (7)-(19).*

It is true that the notation seems inconsistent. We simply switched to the (t, tau) variables for final expressions, and kept (t,t') for the intermediate steps. However, after Eq 33, we use (t,tau) for the calculation of the bath correlation functions so as to emphasize their stationarity property.

• *Something appears to be wrong with equation (19), take for example* $L_1 = A \otimes B$
Then starting from equation (16) one has

$$e^{+\mathcal{L}_S t}\mathrm{Tr}_B\{(e^{-\mathcal{L}_S t}Ae^{+\mathcal{L}_S t}\otimes e^{-\mathcal{L}_B t}Be^{+\mathcal{L}_B t})(e^{-\mathcal{L}_S t'}Ae^{+\mathcal{L}_S t'}\otimes e^{-\mathcal{L}_B t'}Be^{+\mathcal{L}_B t'})\rho_B^e\}e^{-\mathcal{L}_S t}\sigma(t) =$$

$$\mathrm{Tr}_B\{(A\otimes e^{-\mathcal{L}_B t}Be^{+\mathcal{L}_B t})(e^{\mathcal{L}_S(t-t')}Ae^{-\mathcal{L}_S(t-t')}\otimes e^{-\mathcal{L}_B t'}Be^{+\mathcal{L}_B t'})\rho_B^e\}\sigma(t) =$$

$$\mathrm{Tr}_B\{(\tilde{A}(0)\otimes\tilde{B}(0))(\tilde{A}(t-t')\otimes\tilde{B}(t-t'))\rho_B^e\}\sigma(t) =$$

$$\mathrm{Tr}_B\{\tilde{\mathcal{L}}(0)\tilde{\mathcal{L}}(t-t')\rho_B^e\}\sigma(t),$$

Notation inconsistencies in the Schrödinger representation have been corrected and the complete derivation is given in an Appendix.